



## CFD study of Convective Heat Transfer of Water Flow Through Micro-Pipe with Mixed Constant Wall Temperature and Heat Flux Wall Boundary Conditions

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### ABSTRACT

The dissipation of heat in tiny engineering systems can be achieved with fluid flow through micro pipes. They have the advantage of less volume to large surface ratio convective heat transfer. There are deep-rooted analytical relations for convective heat transfer available for fluid flow through macro size pipes. But differences exist between the convective heat transfer for fluid flow through macro and micro pipes. Therefore, there is a good scope of work in micro convection heat transfer to study the mechanism of fundamental flow physics. There have been studies with either constant heat flux wall boundary conditions or constant wall temperature boundary conditions with constant and variable property flows. In this article, first, the numerical simulations are validated with the experimental data for 2D axisymmetric conventional pipe with pipe diameter of 8 mm is taken with laminar, steady, and single-phase water flows with constant wall heat flux boundary condition of 1 W/cm<sup>2</sup>. The computed Nusselt number is compared to the experimental results at different Reynolds numbers of 1350, 1600 and 1700. In the next study, three-dimensional micropipe laminar flow is studied numerically using water with an inlet velocity of 3 m/s and pipe diameter of 100 μm. The mixed wall boundary conditions with upper half pipe surface subjecting to constant wall temperature of 313 K and lower half surface subjecting to 100 W/cm<sup>2</sup> are used in the simulations. The focus of research would be to consider the effect of temperature-dependent properties like thermal conductivity, viscosity, specific heat, and density (a combined effect we call it as variable properties) on micro-pipe flow characteristics like Nusselt number at mixed wall boundary conditions and compare it with the constant property flows. The conventional pipe showed no significant difference with variable and constant property flows with different Reynolds numbers. On contrary the flow through 3D micropipe shows that the Nusselt number with variable property flows is less as compared to the constant property flows.

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## 1. Introduction

The miniature modern machines in medical, chemical and information technology demand highly effective and space-efficient heat exchanging devices. The heat transfer and hydrodynamic features of micro ducts have been investigated enormously in past few years by following researchers: Gulhane and Mahulikar [1], Mala *et al.*, [2], Sobhan and Garimella [3], Morini [4], Hetsroni *et al.*, [5], Kandlikar and Steinke [6], Hong and Asako [7], van Rij *et al.*, [8], Kandlikar [9], Garimella and Sobhan [10], Prabhu and Mahulikar [11], Adham *et al.*, [12], Dixit and Ghosh [13], Rastogi and Mahulikar [14], Tharkar and Mahulikar [15], Tharkar and Mahulikar [16] and Kumar *et al.*, [17]. There are still great variations in the reported data. Alterations exist in many aspects: drag of fluid in laminar, transit and turbulent single-phase flows, heat exchange of fluids, two-phase adiabatic and heated micro-channels. These contexts have difficulties linking to channel shape, aspect ratio, type of fluid, compressibility, phase, surface irregularities and heat transfer circumstances in micro pipes. Different cross sections like triangular and rectangular shapes were taken to study the performance of the hybrid micro-channel flows [18,19]. The traditional heat transport theory tells when the cross-section of channels reduces, the coefficient of heat transfer surges for fully developed fluid flow [20]. Hence, small-sized ducts are selected to get higher values of heat transfer. For instance, these micro conduits are employed to bring about heat loss in Large-Scale Integration (LSI) chips and micro total chemical analyzing system micro Total Analysis System ( $\mu$ -TAS), micro electrical-mechanical systems (MEMS) [21,22].

The liquid flow through micro ducts is used to dissipate heat in miniature engineering systems, which is of paramount importance for their performance. Fluid flow, through conventional macro size ducts, has been extensively studied to develop the well-established analytical relations of heat transfer [5,23,24]. Unfortunately, there exist many discrepancies between the mechanism of fluid flow and heat transfer in macro ducts and micro ducts. This opens a broad door to unveil the flow physics in micro convection heat transfer. The first micro-channel experimental studies performed by Tuckerman and Pease [25], they have demonstrated that a greater heat transfer rate could be achieved with smaller diameters. Subsequently, several investigations have been conducted over the past years to practically analyze the flow through micro-channels [26-29]. However, experimental measurements for micro-channel pipes face a major challenge to date. As an alternative, numerical simulations have been employed by researchers, to overcome this limitation [1,30].

Generally, constant properties ( $C_p$ ) are assumed in fluid flow and heat transfer analytical calculations for duct flows. The effect of fluid property variations with temperature such as density dependent on temperature  $\rho(T)$ , thermal conductivity dependent on temperature  $k(T)$ , specific heat at constant pressure  $C_p(T)$  and viscosity dependent on temperature  $\mu(T)$  are generally neglected that account for large temperature differences when compared with constant properties. It is found that there are around 30% deviations in Nusselt number ( $Nu$ ) [31-34]. Experimental shreds of evidence, with a comprehensive review of literature, were provided by Kandlikar and Steinke [6] showing the discrepancies present in literature for frictional factor and pressure drop for flows through micro-channels. After reviewing of the pertinent literature, it is found that there are suitable discrepancies in the behavior of thermophysical properties as a function of temperature in the micropipe flows with convection.

There are also discrepancies in the comparison with experimental data in pressure drop and heat transfer in micro pipe flows. Hence, there is a need for appropriate modelling of thermophysical properties [ $\rho(T)$ ,  $k(T)$ ,  $C_p(T)$  and viscosity  $\mu(T)$ ] as a function of temperature and their incorporation in the real life applications with the aid of computational results leading to safer designs. The focus of the current work is to investigate the effects of the temperature variation of thermophysical

properties on heat transfer and hydrodynamic characteristics of 3D micropipe with mixed wall boundary conditions. Also, the mixed wall boundary conditions are not studied till date. The objective of the present study is to compute numerically three-dimensional incompressible, steady, laminar flow through micro pipe to investigate the effect of temperature-dependent thermophysical properties on hydrodynamic and heat transfer characteristics at combined constant wall heat flux and constant wall temperature boundary condition. So, as a fundamental fluid dynamic it would be interesting to study the behaviour of heat transfer characteristics i.e., Nu, for developing water flow through the micro pipe with constant wall temperature at top half and constant heat flux applied and bottom half surfaces of a three dimensional (3D) cylindrical pipe. The water is chosen and used as the working fluid in the present computations.

The paper is organized as follows. First, the simulation methodology is discussed. In this section, governing equations, geometry details, grid convergence study, numerical method and formulas are discussed. Next, the CFD validation is performed for conventional pipe flow with available experimental data by Heyhat *et al.*, [35] for the constant property (CP) flows and variable property (VP) flows. Also, the effect of inlet Reynolds number on Nusselt number, skin friction and pressure drop are studied. Next, the 3D cylindrical pipe CFD analysis is done with constant property (CP) flows and variable property (VP) flows. Finally, the conclusions are discussed.

## 2. Numerical Methodology

The 3D form of governing equations in steady-state form for incompressible flows in cylindrical coordinates  $(r, \theta, z)$  are given by Eq. (1)-(5) [36].

Balance of mass:

$$\frac{1}{r} \frac{\partial(ru_r)}{\partial r} + \frac{1}{r} \frac{\partial(ru_\theta)}{\partial \theta} + \frac{\partial u_z}{\partial z} = 0 \quad (1)$$

Balance of momentum:

$$\rho \left( \frac{\partial u_r}{\partial r} + \frac{u_\theta}{r} \frac{\partial u_r}{\partial \theta} - \frac{u_\theta^2}{r} + u_z \frac{\partial u_r}{\partial z} \right) = -\frac{\partial p}{\partial r} + \mu \left( \frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial}{\partial r} (ru_r) \right) + \frac{1}{r^2} \frac{\partial^2 u_r}{\partial \theta^2} - \frac{2}{r^2} \frac{\partial u_\theta}{\partial \theta} + \frac{\partial^2 u_r}{\partial z^2} \right) \quad (2)$$

$$\rho \left( \frac{\partial u_\theta}{\partial r} + \frac{u_\theta}{r} \frac{\partial u_\theta}{\partial \theta} + \frac{u_r u_\theta}{r} + u_z \frac{\partial u_\theta}{\partial z} \right) = -\frac{1}{r} \frac{\partial p}{\partial \theta} + \mu \left( \frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial}{\partial r} (ru_\theta) \right) + \frac{1}{r^2} \frac{\partial^2 u_\theta}{\partial \theta^2} + \frac{2}{r^2} \frac{\partial u_r}{\partial \theta} + \frac{\partial^2 u_\theta}{\partial z^2} \right) \quad (3)$$

$$\rho \left( \frac{\partial u_z}{\partial r} + \frac{u_\theta}{r} \frac{\partial u_z}{\partial \theta} + u_z \frac{\partial u_z}{\partial z} \right) = -\frac{\partial p}{\partial z} + \mu \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_z}{\partial \theta^2} + \frac{\partial^2 u_z}{\partial z^2} \right) \quad (4)$$

Balance of energy:

$$\rho C_p \left( u_r \frac{\partial T}{\partial r} + \frac{u_\theta}{r} \frac{\partial T}{\partial \theta} + u_z \frac{\partial T}{\partial z} \right) = k \left[ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 T}{\partial \theta^2} + \frac{\partial^2 T}{\partial z^2} \right] + 2\mu \left\{ \left( \frac{\partial u_r}{\partial r} \right)^2 + \left[ \frac{1}{r} \left( \frac{\partial u_\theta}{\partial \theta} + u_r \right) \right]^2 + \left( \frac{\partial u_z}{\partial z} \right)^2 \right\} + \mu \left\{ \left( \frac{\partial u_\theta}{\partial z} + \frac{1}{r} \frac{\partial u_z}{\partial \theta} \right)^2 + \left( \frac{\partial u_z}{\partial r} + \frac{\partial u_r}{\partial z} \right)^2 + \left[ \frac{1}{r} \frac{\partial u_r}{\partial \theta} + r \frac{\partial}{\partial r} \left( \frac{u_\theta}{r} \right) \right]^2 \right\} \quad (5)$$

## 2.1 Modelling of Thermo Physical Properties for Water

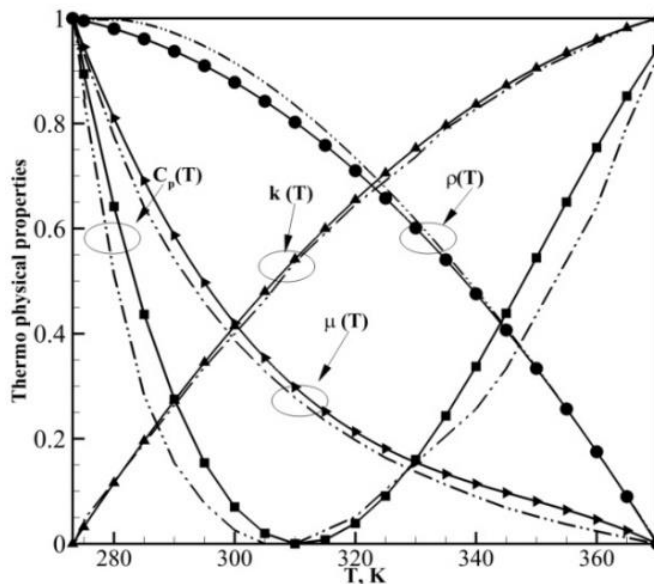
These polynomial functions of thermophysical properties are used in our numerical simulations. The variation of thermophysical properties is shown in Figure 1. It shows that  $\rho(T)$  and  $\mu(T)$  decreases  $k(T)$  increases whereas  $C_p(T)$  decreases up to 310 K and increases beyond it, up to 373.15 K, with increasing temperature. The modeling of thermophysical properties as a function of temperature is given by,

$$\rho(T) = 765.33 + 1.8142(T) - 0.0035(T)^2 \quad (6)$$

$$C_p(T) = 1.095e^4 - 59.27(T) + 0.171(T)^2 - 0.0001623(T)^3 \quad (7)$$

$$k(T) = -0.5752 + 6.3967 \times 10^{-3}(T) - 8.151 \times 10^{-6}(T)^2 \quad (8)$$

$$\mu(T) = 9.67 \times 10^{-2} - 8.207 \times 10^{-4}(T) + 2.344 \times 10^{-6}(T)^2 - 2.244 \times 10^{-9}(T)^3 \quad (9)$$



**Fig. 1.** Assessment of property polynomial functions to the physical data with the variation of temperature for water [37]

## 2.2 CFD Validation with Experiments

The grid is generated using the Pointwise package and the ANSYS-Fluent package 19.3 is used to carry out the numerical simulations. The post-processing is done using the Tecplot package. The governing equations are discretized using the second-order finite volume technique. The SIMPLE scheme is used to couple the velocity and pressure variables. The implicit method is used to reach the steady-state solutions. The no-slip velocity and uniformly distributed constant heat flux ( $q_w$ ) boundary conditions are used at the wall. The velocity inlet boundary conditions are given in Table 1. Atmospheric pressure is assigned at the outlet of pipe,  $p_{out} = 1.01325 \times 10^5$  Pa. An axisymmetric boundary condition i.e., gradients in an axial direction,  $\partial(\rho)/\partial z$ ,  $\partial(T)/\partial z$  and  $\partial(\rho)/\partial z$  are assumed to be zero at the centerline axis.

**Table 1**  
Geometry and boundary conditions for different test cases

Parameters	Conditions
Pipe diameter (D), m	0.008
Inlet temperature ( $T_{in}$ ), K	283
Inlet velocities ( $u_{in}$ ), $ms^{-1}$	0.02784, 0.0262, 0.0221
Reynolds number (Re)	1700, 1600, 1350
Heat Flux ( $q_w$ ), $W/m^2$	10610

The Nusselt number for constant properties ( $Nu_{CP}$ ) and variable properties ( $Nu_{VP}$ ) for constant wall heat flux (CHF) are numerically calculated by Eq. (10) and Eq. (11),

$$Nu_{CP} = \frac{hD}{k_{CP}} = \frac{q_w D}{k_{CP} (T_w - T_m)_{CP}} \quad (10)$$

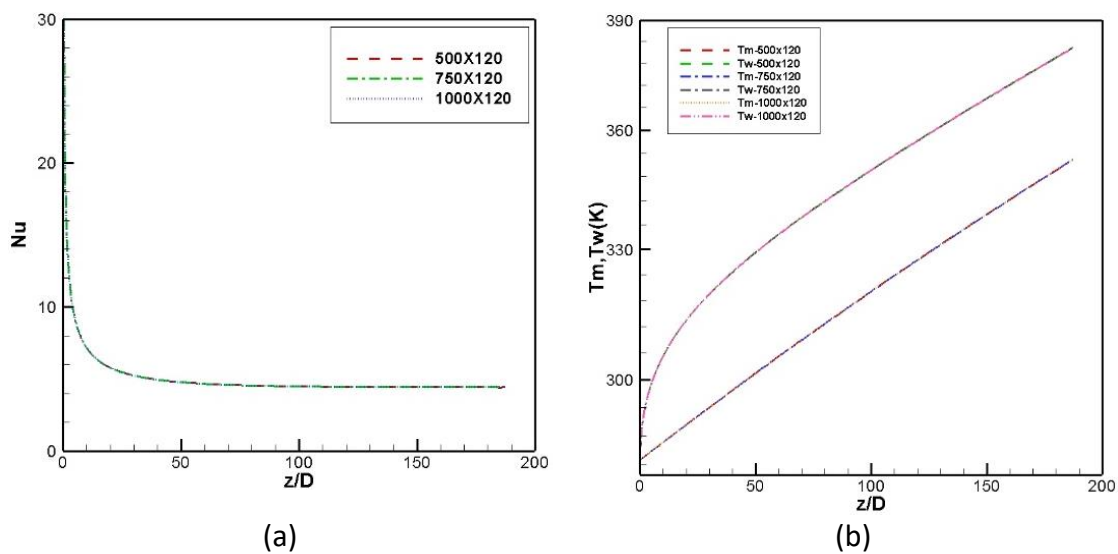
$$Nu_{VP} = \frac{hD}{k(T)} = \frac{q_w D}{k(T) (T_w - T_m)_{VP}} \quad (11)$$

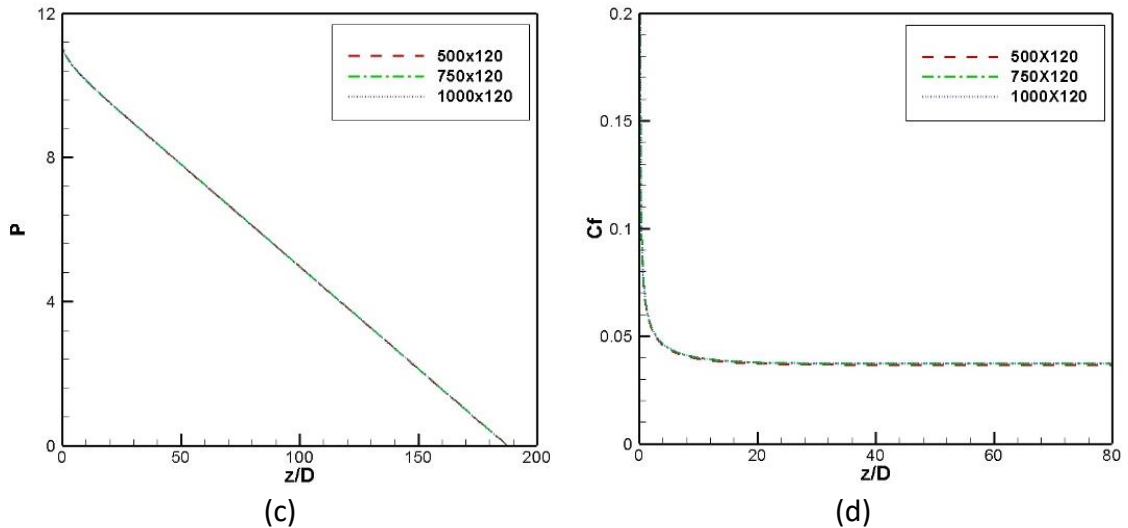
Here, (h) denotes the heat transfer coefficient, the ( $q_w$ ) is constant wall heat flux.

In the numerical simulations, the conventional pipe length (L) of 1.5 m and its diameter (D) as 0.008 m. The mesh is exponentially stretched in the radial direction from the wall and the first cell near the wall is taken as  $5 \times 10^{-6}$  m. The grid is also exponentially stretched in the z-direction with many grid points at the entrance region to capture high gradients of the developing flow. The grid convergence analysis is done in detail.

### 2.3 Grid Convergence Study

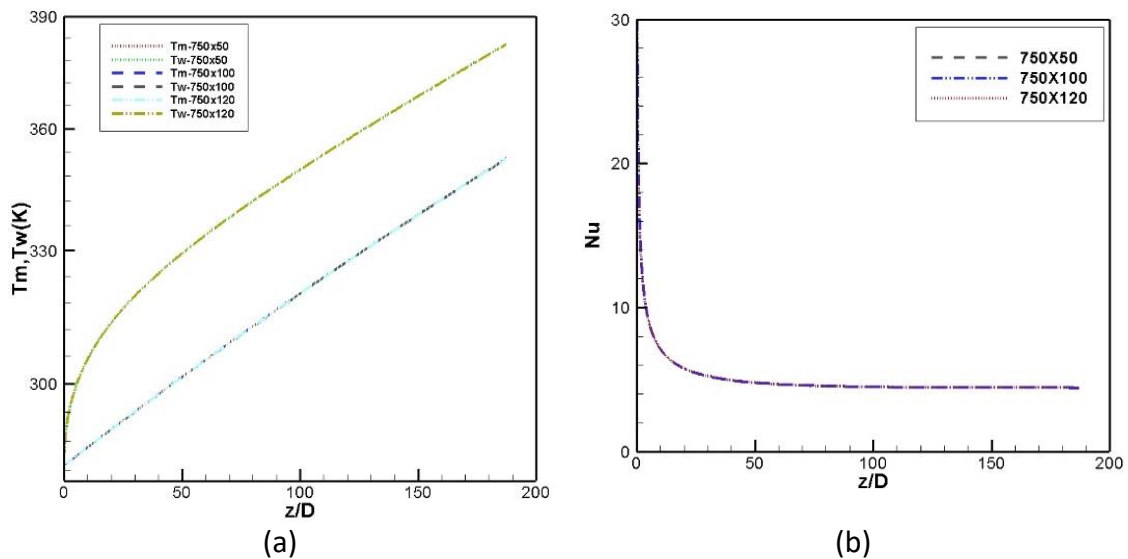
Figure 2 shows that increase in number of grid points in the axial direction from 500 to 1000 (while keeping same number of grid points of 120 in radial direction) has minor effect on the  $T_m$ ,  $T_w$ , Nu, P and  $C_f$  properties and hence 750x120 is taken as converged grid i.e., the solution does not change with further increase in number of grid points from 750 to 1000.

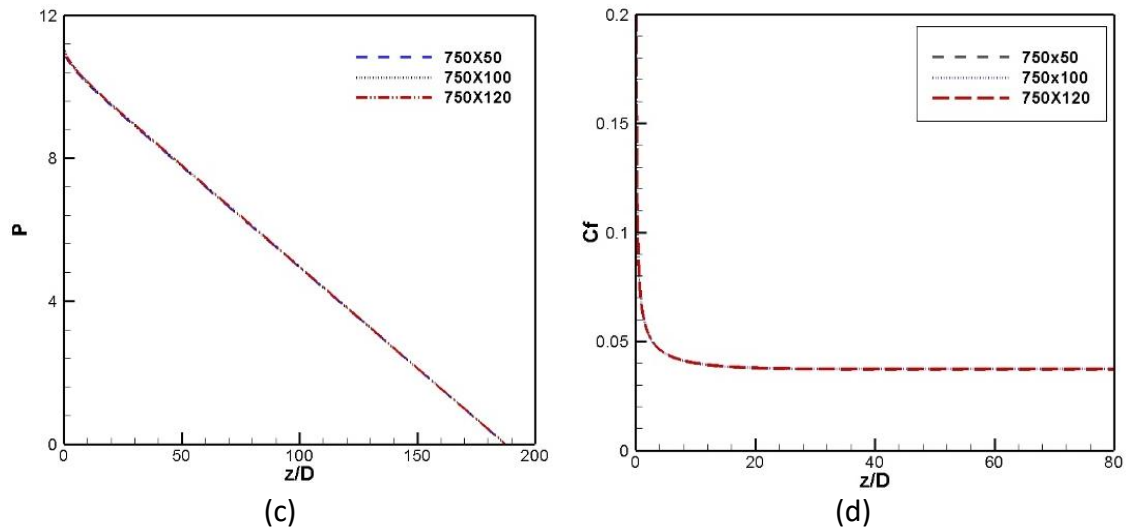




**Fig. 2.** Variation of (a) mean temperature  $T_m$ , wall temperature  $T_w$  (b) Nusselt number (c) pressure and (d) skin friction with different grids of 500x120, 750x120 and 1000x120 with increase in number of points in the axial direction for  $Re = 1700$  at constant wall heat flux =  $10610 \text{ W/m}^2$

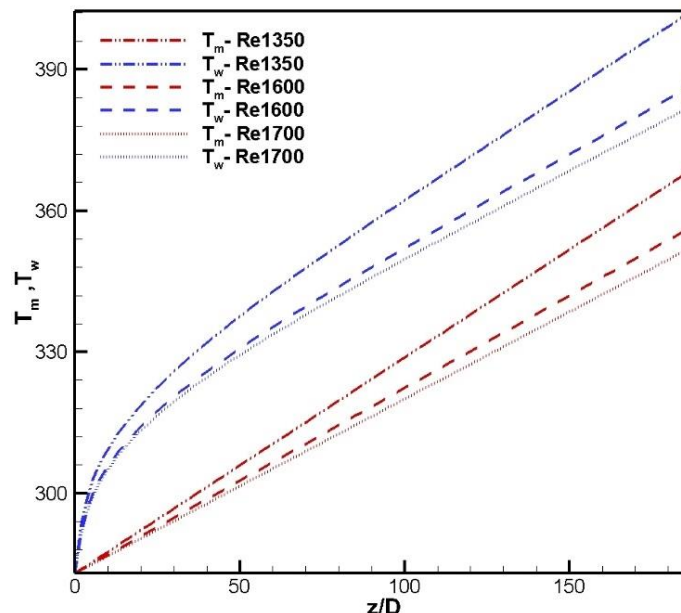
Figure 3 shows that that increasing the number of points in the radial direction from 50 to 100 and then to 120 (while keeping same number of grid points of 750 in the axial direction) has a small change in  $T_m$ ,  $T_w$ ,  $Nu$ ,  $P$  and  $C_f$ . Therefore, 750 points in the axial direction and 100 points in the radial direction are sufficient to capture the velocity and temperature gradients. We conclude that the  $750 \times 100$  grid is taken as a converged grid i.e, the solution does not alter with further increase in both the axial and radial directions and hence this grid is used in two dimensional (2D) Computational Fluid Dynamics (CFD) calculations for conventional pipe.



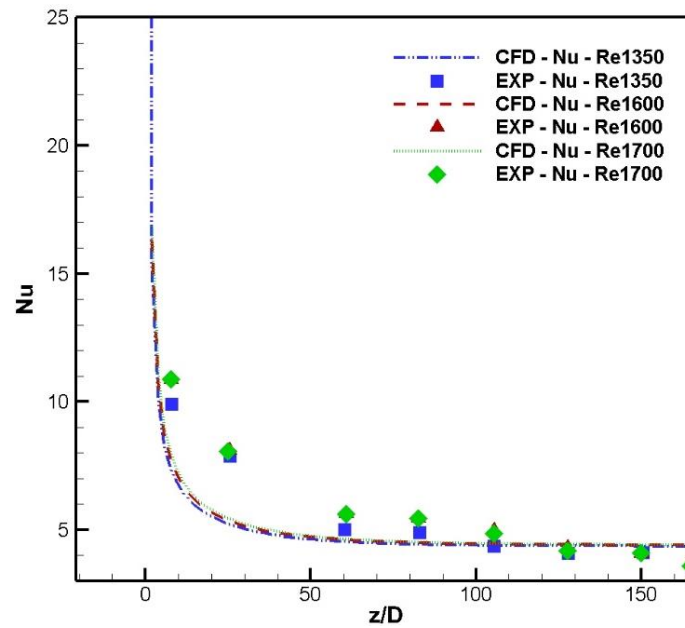


**Fig. 3.** Variation of (a) mean temperature  $T_m$ , wall temperature  $T_w$ , (b) Nusselt number (c) pressure and (d) skin friction with different grids of 750x50, 750x100 and 750x120 with increase in number of points in the radial direction for  $Re = 1700$  at constant wall heat flux =  $10610 \text{ W/m}^2$

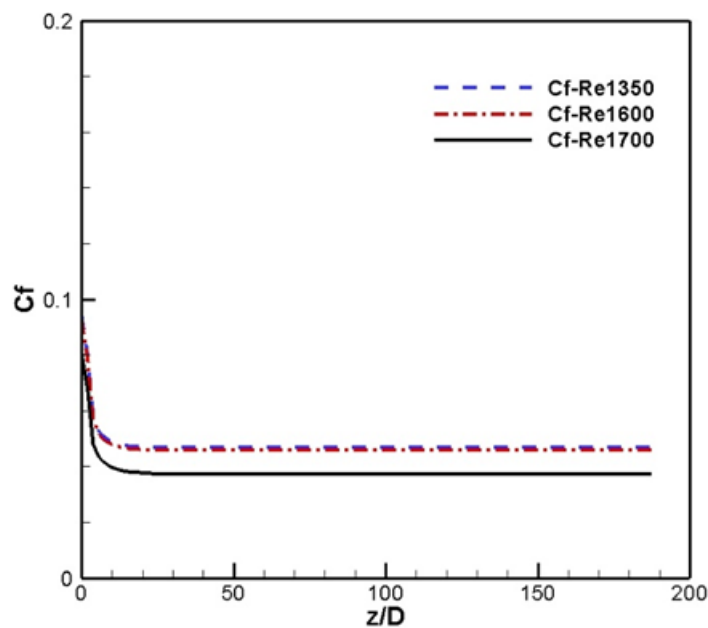
Figure 4 shows that the variation of mean and wall temperatures with constant property flow with constant heat flux boundary condition for three Reynolds numbers of 1350, 1600 and 1700, a constant temperature difference  $[\Delta T = (T_w - T_m)]$  is observed beyond  $(z/d) = 50$ , therefore depicting thermally fully developed flow. The variation of Nusselt number along the axial direction under similar conditions of Reynolds number and constant wall heat flux is computed and compared with experimental data of Heyhat *et al.*, [35] as shown in Figure 5. The variation of skin friction factor,  $C_f$  along the walls of the micropipe is shown in Figure 6 shows that the  $C_f$  decreases with an increase in the Reynolds number because of the corresponding increase in the boundary layer thickness.



**Fig. 4.** Variation of mean temperature ( $T_m$ ) and wall temperature ( $T_w$ ) long the axial direction for the constant properties CP for  $Re = 1350, 1600$  and  $1700$  at constant wall heat flux =  $10610 \text{ W/m}^2$



**Fig. 5.** Variation of Nusselt number along the axial direction for the constant properties CP for Re = 1350, 1600 and 1700 at constant wall heat flux = 10610 W/m<sup>2</sup> compared to the experimental data by Heyhat *et al.*, [35]



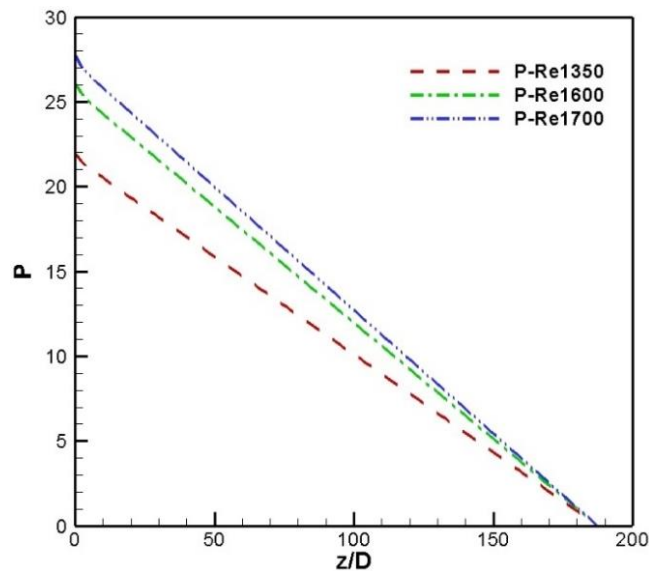
**Fig. 6.** Variation of skin friction ( $C_f$ ) along the wall for the constant properties CP for Re = 1350, 1600 and 1700 at constant wall heat flux  $q_w = 10610$  W/m<sup>2</sup>

Figure 7 shows the variation of pressure along the micro-pipe walls for different Reynolds numbers and constant wall heat flux boundary conditions. The low Reynolds number has a lower pressure drop while the higher Reynolds number has a higher which is as per expectation. In the present work, we have validated our CFD simulation data with the corresponding experimental results. The comparison between CFD and experimental data for Nusselt number with Reynolds number is shown in Figure 8. There is an excellent agreement between the CFD simulation results

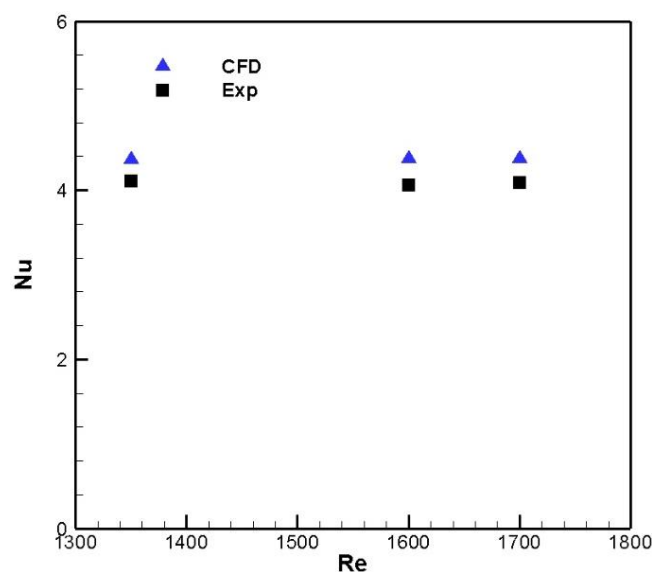


and the experimental data. These small deviations may be attributed due to the uncertainty in the experimental measurements.

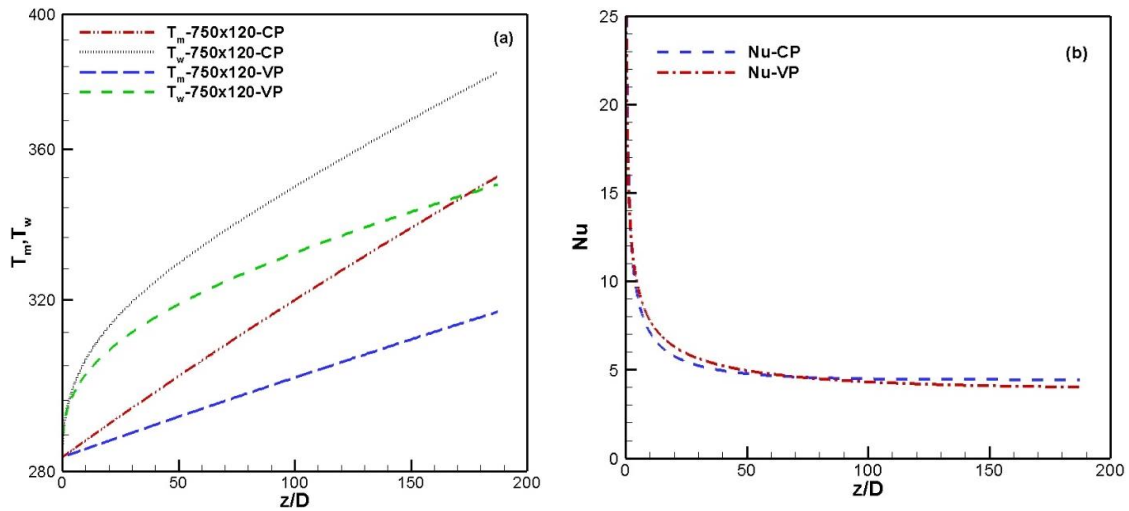
In case of macro-pipes, the simulation results show a peculiar variation between the mean bulk fluid temperatures and wall temperatures for variable properties shown in Figure 9(a). This may be attributed to the presence of strong thermal gradients which influence the thermophysical properties such as density  $\rho(T)$ , thermal conductivity  $k(T)$ , specific heat at constant pressure  $C_p(T)$  and dynamic viscosity  $\mu(T)$ . The effect of variable properties on of Nusselt number is shown in Figure 9(b). The comparison of Nusselt number with non-dimensional length along the flow direction of macro-pipe reveals that overall, the influence of variable properties is not having a significant effect. But the effects of the thermal gradients near-wall influence the thermophysical properties in micro pipes. In the next section, these results are discussed.



**Fig. 7.** Variation of wall pressure along the pipe with constant properties CP for Re = 1350, 1600 and 1700 at constant wall heat flux =  $10610 \text{ W/m}^2$



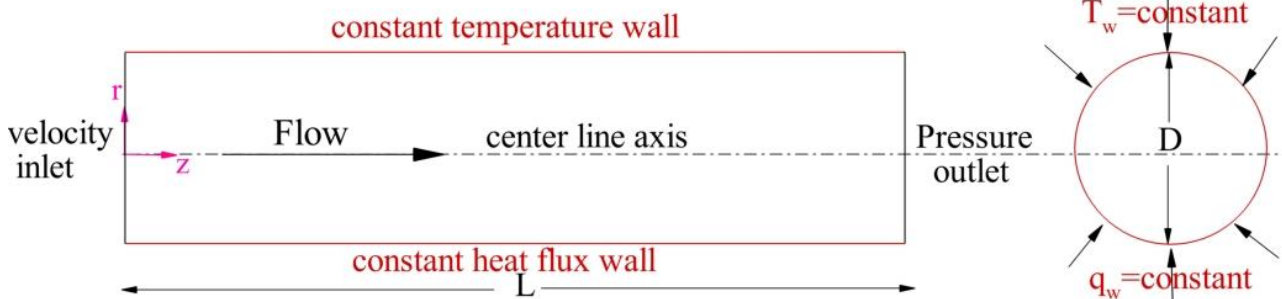
**Fig. 8.** Comparison of CFD and experimental data for variation of Nusselt number with Reynolds number



**Fig. 9.** Variation of (a) mean temperature  $T_m$  and wall temperature  $T_w$  and (b) Nusselt number with constant properties (CP) and variable properties (VP) for flow in a pipe (750 x120 grid points) with  $Re = 1700$  at constant wall heat flux =  $10610 \text{ W/m}^2$

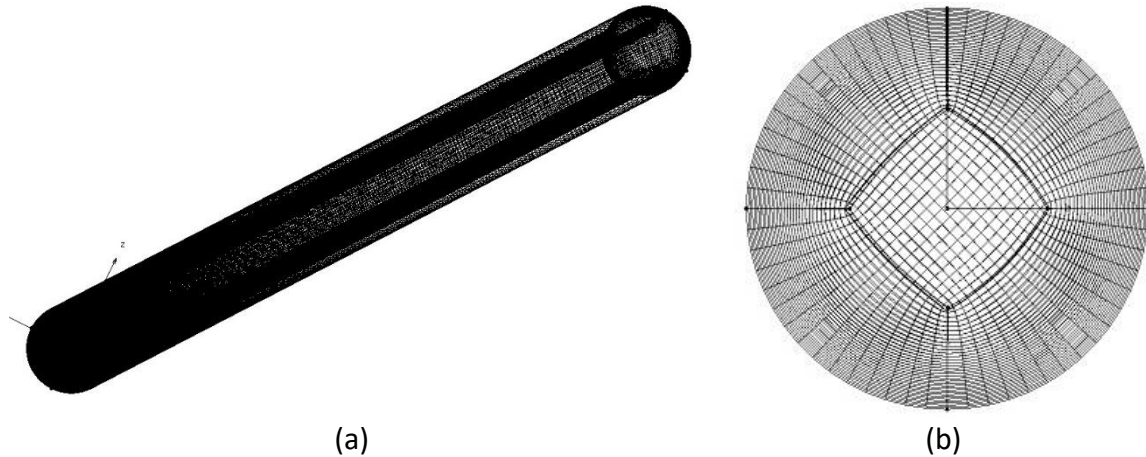
### 3. Results and Discussions

In this section, we discuss the three-dimensional CFD results for micro-pipe water flow as shown in Figure 10 with an inlet velocity of  $3 \text{ ms}^{-1}$  and a pipe diameter of  $100 \times 10^{-6} \text{ m}$ . The mixed wall boundary conditions are (a) upper half pipe - constant wall temperature =  $313.15 \text{ K}$  and (b) lower half pipe - constant heat flux =  $100 \text{ W/cm}^2$ .



**Fig. 10.** Micropipe convective heat transfer flow with combined constant heat flux and constant wall temperature boundary condition

Figure 11 shows the 3D pipe mesh and its front view is also shown (H-topology grid). Care is taken to remove the singularity developed when we use rotate option in the grid generation package to generate the grid from a 2D rectangle to form a 3D pipe mesh. A grid converged solution is obtained on  $300 \times 40 \times 40$  mesh with 300 points in the  $x$ -direction, 40 points in the  $y$ -direction and 40 points in the  $z$ -direction.



**Fig. 11.** Meshing views (a) A micro pipe 3D mesh and it's (b) Front view

The formulas used for convective heat transfer equations at mixed boundary conditions of constant heat flux and constant wall temperature are discussed as follows.

The  $q_w$  is constant wall heat flux at the bottom surface with  $100 \text{ W/cm}^2$  and, we compute values of  $T_w$  at the bottom surface and the mean temperature,  $T_m$  ( $T_m$ )<sub>CP</sub> and ( $T_m$ )<sub>VP</sub> for constant property (CP) and variable property flows (VP) separately from the Fluent CFD post package. Next, the Nusselt number at CP and VP flows are calculated using Eq. (10) and Eq. (11).

The Nusselt number at CP and VP flows with upper half wall constant wall temperature (CWT) equal to  $(T_w)_{CP} = (T_w)_{VP} = 373.15 \text{ K}$  and taking are calculated by Eq. (12) and Eq. (13). The mean temperature,  $(T_m)_{CP}$  and  $(T_m)_{VP}$  are calculated from the Fluent CFD post package in below equations. Here,  $q_w$  increases from pipe inlet to outlet at the upper half surface. It is calculated at different locations along the centerline of the upper semi-cylindrical wall from the Fluent CFD post package.

$$\{Nu_{CP}\}_{CWT} = \frac{hD}{k_{CP}} = \frac{q_w D}{k_{CP} (T_w - T_m)_{CP}} \quad (12)$$

$$\{Nu_{VP}\}_{CWT} = \frac{hD}{k_{VP}} = \frac{q_w D}{k_{VP} (T_w - T_m)_{VP}} \quad (13)$$

For the mixed boundary condition of CHF and CWT we take the weights and calculate Nu as follows. From theory, the  $Nu = 4.36$  is taken for hydrodynamic and thermally developed flow for CHF and the  $Nu = 3.66$  is taken for hydrodynamic and thermally developed flow for CWT and their weights are calculated as,  $CHF = 4.36 / (4.36 + 3.66) = 54.4\%$ , which gives weight for  $CHF_{wt} = 0.544$  and therefore weight for  $CWT_{wt} = (1 - 0.544) = 0.456$ . Therefore, the Nusselt number for mixed wall boundary condition as shown in Figure 8 for CP flows is given by,

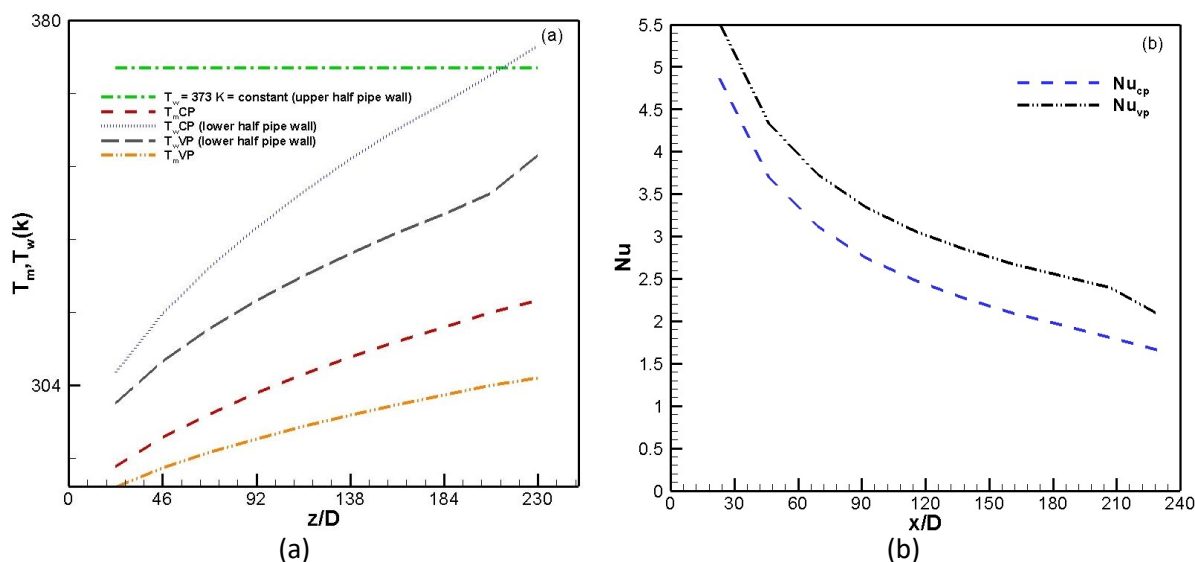
$$\{Nu_{CP}\}_{(CHF + CWT)} = CHF_{wt} \times \{Nu_{CP}\}_{CHF} + CWT_{wt} \{Nu_{CP}\}_{CWT} \quad (14)$$

Similarly, the Nu for mixed wall boundary condition for VP flows is given by,

$$\{Nu_{VP}\}_{(CHF + CWT)} = CHF_{wt} \{Nu_{VP}\}_{CHF} + CWT_{wt} \{Nu_{VP}\}_{CWT} \quad (15)$$

Figure 12(a) shows the variation of bulk mean temperature ( $T_m$ ) and micro-pipe wall temperature ( $T_w$ ) along the flow direction. The driving temperature difference for constant properties (CP) is higher than the variable properties (VP) under the specialized boundary conditions acting on pipe upper and lower halves. Figure 12(b) shows the variation of the Nusselt number along with the

dimensionless distance of the micro-pipe. The influence of variable properties (thermo-physical) on heat transfer is being captured more accurately. We mean to say that the calculations of the Nusselt number based on the constant properties (CP) have shown to under (lower) prediction of heat transfer. This implies that the design of micro-pipes using variable properties (VP) is a safer design under the special boundary conditions of constant heat flux on one half and constant wall temperature on the other half of micro pipe walls.



**Fig. 12.** Variation of computed (a) mean temperature  $T_m$  and wall temperature  $T_w$  and (b) Nusselt number with constant properties CP and variable properties VP for flow in a 3D pipe with an inlet velocity of  $3\text{ ms}^{-1}$  with mixed boundary conditions: at constant wall temperature of  $313.15\text{ K}$  (top half-pipe surface) and constant wall heat flux of  $100\text{ W/cm}^2$  at the bottom half surface pipe

#### 4. Conclusions

The CFD analysis is performed on the water flowing through the 2D conventional pipe subjected to constant wall heat flux BC and 3D micro-pipe with mixed constant wall temperature and constant heat flux BC. The study concludes that:

- i. The 2D pipe flow CFD results of the Nusselt number matches the experimental data. The Nusselt number does not alter whereas the skin-friction decreases, and the pressure drop increases with an increase in inlet Reynolds number. Though there are differences in the variation of system-dependent parameters  $T_m$  and  $T_w$  along the length of pipe for CP and VP flows. The variation of system-independent parameters Nu with Re does not vary because of fewer gradients near the wall for the conventional pipe.
- ii. The understanding of the influence of variable properties (VP) on heat transfer enhancement in micro-pipe geometry is significant and this will aid us for efficient, safer designs of micro-pipes under specialized boundary conditions of both constant wall temperature and constant heat flux operating simultaneously.

The effect of inlet Reynolds number (Re) on Nusselt number (Nu) and friction factor will be studied for these fluid flows in the future.

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