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# MHD and Slip Effect in Micropolar Hybrid Nanofluid and Heat Transfer over a Stretching Sheet with Thermal Radiation and Non-uniform Heat Source/Sink

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ARTICLE INFO	ABSTRACT
Article history: Received 23 September 2020 Received in revised form 21 November 2020 Accepted 25 November 2020 Available online 30 November 2020	In the presence of slips, non-uniform heat source/sink, thermal radiation and magnetohydrodynamic (MHD), micropolar hybrid nanofluid and heat transfer over a stretching sheet has been studied. The problem is modelled as a mathematical formulation that involves a system of the partial differential equation. The similarity approach is adopted, and self-similar ordinary differential equations are obtained and then those are solved numerically using the shooting method. The flow field is affected by the presence of physical parameters such as micropolar parameter, magnetic field parameter, suction parameter and slip parameter whereas the temperature field is affected by thermal radiation parameter, space-dependent parameter, temperature-dependent internal heat generation/absorption parameter, Prantl number and Biot number. The skin friction coefficient, couple stress and local Nusselt number are tabulated and analysed. The effects of the governing parameters on the velocity profiles, angular velocity profiles and temperature profiles are illustrated graphically. The results of velocity profiles, angular velocity profiles and temperature profiles are also obtained for several values of each parameters involved.
Hybrid nanofluid; micropolar fluid; thermal radiation; suction;	
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## 1. Introduction

Convective heat transfer has played as one of the biggest important roles in many applications especially in engineering field, technologies, medical sciences and natural processes. Heat transfer in conventional fluids such as ethylene glycol, water and oil are very important in our industrial. However, to enhance the performance of heat transfer, hybrid nanofluid has been introduced as a new class of nanofluid, marked by its thermal properties and potential utilities which serve the

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purpose. The main agenda of the current analysis is to present the behavior of emerging hybrid nanofluid in the presence of micropolar fluid theory and rotation over a stretching sheet. Eringen [1] was first introduced the micropolar theory and discussed it in detail in his research. A lot of researchers (see Ahmadi [2], El-Hakiem [3], Damseh *et al.*, [4], Ishak [5], Das [6] and Naveed *et al.*, [7]) started to studied and did an intensive investigation by extending Eringen's work by considering many aspects.

However, many other researchers have shift their interest and come out with different brilliant investigation with various type of hybrid nanoparticles using a new procreation of heat transfer fluid, which is hybrid nanofluid to give a better heat transfer outcome compared to traditional cooling fluid. Devi and Devi [8] were the first to analyse the boundary layer equations of hybrid nanofluid by introducing a new special form of thermophysical properties.

Suresh *et al.*, [9] conducted a fully developed laminar convective heat transfer and pressure drop characteristics through a uniformly heated circular tube using  $Al_2O_3 - Cu$  /water hybrid nanofluid. Besides, Momin [10] investigated the mixed convection laminar flow of hybrid nanofluid  $Al_2O_3 - Cu$ /water inside an inclined tube. Furthermore, there are many other researchers also focusing on different problems by hybrid different nanoparticles with different effects and fluids. Moreover, Subhani and Nadeem [11] conducted a numerical analysis of micropolar hybrid nanofluid and the outcome proved that the rate of heat transfer of hybrid nanofluid is higher than nanofluid even with various effects of parameter such as micropolar effect, rotation and porosity. Khashi'ie *et al.*, [12] were focusing their studies on mixed convection stagnation point flow of hybrid nanofluid over a permeable stretching/shrinking cylinder with the presence of thermal stratification effect.

Inspired by all the previous literatures, the motivation of this studies is to investigate the effects of the slip-on hybrid micropolar nanofluid flow and heat transfer over a stretching sheet with thermal radiation, non-uniform heat source/sink and magnetohydrodynamic (MHD). The nonlinear ordinary differential equations obtained are solved numerically by shooting method. The effects of governing parameters on the skin friction coefficients, wall couple stress and local Nusselt number are investigated and analysed. We extended the Pal and Mandal [13] studies by adding the effect of the slip and MHD to their problem.

## 2. Methodology

We consider the steady two-dimensional flow of micropolar hybrid nanofluid past a stretching sheet in the presence of thermal radiation with slip, non-uniform heat source/sink and MHD as shown in Figure 1. The fluid velocities are denoted by u and v which are measured in the x – and y – axes, respectively and magnetic field of intensity  $B_0$  is along y direction.



**Fig. 1.** Physical model with coordinate system for stretching sheet



The steady state boundary layer equations for the propose problem in the Cartesian coordinates are given by:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$\rho_{hnf} \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = \left( \mu_{hnf} + \kappa_1 \right) \frac{\partial^2 u}{\partial y^2} + \kappa_1 \frac{\partial N}{\partial y} - \sigma_{hnf} B_0^2 u, \tag{2}$$

$$\rho_{hnf} j \left( u \frac{\partial N}{\partial x} + v \frac{\partial N}{\partial y} \right) = -\kappa_1 \left( 2N + \frac{\partial u}{\partial y} \right) + \gamma_{hnf} \left( \frac{\partial^2 N}{\partial y^2} \right), \tag{3}$$

$$u\frac{\partial T}{\partial x} + \frac{\partial T}{\partial y} = \alpha_{hnf}\frac{d^2T}{dy^2} - \frac{1}{(\rho C_P)_{hnf}}\frac{\partial q_r}{\partial y} - \frac{1}{(\rho C_p)_{hnf}}q^{\prime\prime\prime},\tag{4}$$

subject to the following boundary conditions as given by (Ibrahim and Gadisa [14]):

$$u = u_w(x) + L\frac{\partial u}{\partial y}, v = v_w, N = -m\frac{\partial u}{\partial y}, -k\frac{\partial T}{\partial y} = h_1(T_f - T) \text{ at } y = 0,$$
(5)

$$u \to 0, N \to 0, T \to T_{\infty} \text{ as } y \to \infty.$$
 (6)

Here, T is the fluid temperature, N is the microrotation vector normal to xy-plane, j is the microinertia density,  $\kappa_1$  is the vortex viscosity(gyro-viscosity),  $B_0$  is the strength of magnetic field,  $(\rho C_p)_{hnf}$  is the heat capacitance in the hybrid nanofluid,  $\mu_{hnf}$  is the dynamic viscosity of the hybrid nanofluid,  $\rho_{hnf}$  is the density of the hybrid nanofluid,  $\alpha_{hnf}$  is the thermal diffusivity of the hybrid nanofluid,  $\sigma_{hnf}$  is the electrical conductivity in hybrid nanofluid,  $\gamma_{hnf}$  is the spin gradient viscosity in hybrid nanofluid,  $v_w$  is the constant flux velocity with  $v_w < 0$  for suction and  $v_w > 0$  for injection,  $T_f$  is the convective fluid temperature and  $h_1$  is the heat coefficient. It is assumed that the temperature at the stretching sheet takes the constant values  $T_w$ , while the temperature of the ambient nanofluid take the constant value  $T_\infty$ . These variables are mathematically defined as (Subhani and Nadeem [11]):

$$(\rho C_p)_{hnf} = \phi_2 (\rho C_p)_{s_2} + (1 - \phi_2) \left[ (1 - \phi_1) (\rho C_p)_f + \phi_1 (\rho C_p)_{s_1} \right], \mu_{hnf} = \frac{\mu_f}{(1 - \phi_1)^{2.5} (1 - \phi_2)^{2.5}}, \rho_{hnf} = \phi_2 \rho_{s2} + (1 - \phi_2) \left[ (1 - \phi_1) \rho_f + \phi_1 \rho_s \right], \sigma_{hnf} = \phi_2 \sigma_{s2} + (1 - \phi_2) \left[ (1 - \phi_1) \sigma_f + \phi_1 \sigma_s \right], \frac{k_{hnf}}{k_{nf}} = \frac{(k_{s_2} + 2k_{nf}) - 2\phi_2(k_{nf} - k_{s_2})}{(k_{s_2} + 2k_{nf}) + \phi_2(k_{nf} - k_{s_2})}, \frac{k_{nf}}{k} = \frac{(k_s + 2k_f) - 2\phi_2(k_f - k_1)}{(k_{s_1} + 2k_f) + \phi_1(k_f - k_{s_1})}, \alpha_{hnf} = \frac{k_{hnf}}{(\rho C_p)}$$

$$(7)$$

We note that m is a constant such that  $0 \le m \le 1$ . It is noted that m = 0 indicates N = 0 which is no-spin condition. The case m = 1/2 represents the weak concentration of microelements. The case corresponding to m = 1 is used for the modelling of turbulent flows. Now regarding the approximation of Rosseland for radiation according to Bhattacharyya *et al.*, [15], the radiative heat flux is simplified as  $q_r = (4\sigma^*/3k^*)\partial T^4/\partial y$ , where  $\sigma^*$  is the Stefan–Boltzman constant,  $k^*$  is the absorption coefficient.

The q''' is the non-uniform heat source/sink modelled as (Pal and Chatterjee [16]):



$$q^{\prime\prime\prime} = \left(\frac{\kappa_{hnf} u_w(x)}{x v_f}\right) [A^* (T_w - T_\infty) F'(\eta) + (T - T_\infty) B^*$$
(8)

where  $A^*$  and  $B^*$  represent the space dependant and temperature dependent internal generation or absorption parameters respectively. If the case  $A^* < 0$  and  $B^* < 0$ , it shows that the case corresponds to internal heat absorption otherwise if the case  $A^* > 0$  and  $B^* > 0$ , it shows that the case corresponds to internal heat generation.

The following dimensionless variables are introduced as follows:

$$\eta = (cv)^{\frac{1}{2}}y, \psi = (cv)^{\frac{1}{2}}xf(\eta), N = cx\left(\frac{c}{v}\right)^{\frac{1}{2}}h(\eta), \theta(\eta) = \frac{T - T_{\infty}}{T_w - T_{\infty}},$$
(9)

where  $\eta$  is the similarity variable and  $\psi$  is the stream function defined in the usual way to identically satisfied Eq. (1) as  $u = \frac{\partial \psi}{\partial y}$  and  $v = -\frac{\partial \psi}{\partial x}$ .

By substituting Eqs. (5) - (8) into Eqs. (2) - (4) along with boundary conditions Eqs. (5) and (6) are then transform into the following nonlinear ordinary differential equations obtained as

$$\frac{1}{\phi_2}(\phi_1 + K)f''' - \frac{\phi_4}{\phi_2}Mf' + ff'' - f'^2 + \frac{1}{\phi_2}Kh' = 0$$
(10)

$$\frac{K}{\phi_2}(2h+f'') - \frac{1}{\phi_2}\left(-\phi_1 - \frac{K}{2}\right)h'' + fh' - f'h = 0$$
(11)

$$\frac{k_{hnf}}{Prk_f} + \frac{Nr}{Pr} \{1 - (1 + \theta_w)\theta\}^3 \theta'' - \frac{1}{Pr} (A^*f' + B^*\theta) + \phi_3 f\theta' = 0$$
(12)

subject to boundary conditions

$$f(\eta) = S, f'(\eta) = 1 + \alpha f''(0), h(\eta) = -mf''(\eta), \theta(\eta) = \left(\frac{\theta'(0)}{\gamma} + 1\right) \text{ at } \eta = 0,$$
(13)

$$f'(\eta) \to 0, h(\eta) \to 0, \theta(\eta) \to 0 \text{ as } \eta \to \infty,$$
 (14)

where differentiation with respect to  $\eta$  is signified by ('),  $Pr = \frac{v_f}{a_f}$  is the Prandtl number and  $K = \frac{k_1}{\mu_f}$  is the material/micropolar parameter,  $S = -v_w/(av)^{1/2}$  is the constant flux parameter with S > 0 corresponds to suction and S < 0 corresponds to injection,  $M = \frac{\sigma_f}{c\rho_f}B_0^2$  is the magnetic field,  $Nr = -\frac{16\sigma^*T_\infty^3}{3\kappa_f K^*}$  is thermal radiation parameter,  $\alpha$  is velocity slip parameter and  $\gamma$  is Biot number. We also have (see Pal and Mandal [16]):

$$\theta_{1} = (1-\theta)^{2.5}, \theta_{2} = 1-\phi + \phi\left(\frac{\rho_{s}}{\rho_{f}}\right), \theta_{3} = 1-\phi + \frac{(\rho C_{P})_{s}}{(\rho C_{P})_{f}},$$
  
$$\theta_{4} = 1-\phi + \phi\left(\frac{\sigma_{s}}{\sigma_{f}}\right).$$
(15)

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# 3. Results and Discussion

The nonlinear ordinary differential equations, Eqs. (10)–(12) which are subjected to the boundary conditions Eqs. (13) and (14) are solved numerically using shooting method. Numerical solutions for the effects of partial slip on boundary layer flow in micropolar fluid in the presence of thermal radiation, non-uniform heat source/sink and MHD are reported in this section. The values of the skin-friction coefficient, wall couple stress coefficient and local Nusselt number for various governing parameters are obtained. Table 1 shows the thermophysical properties of base fluid and nanoparticles. We have considered water as the fluid base and the nanoparticle are Cu and  $Al_2O_3$ . Table 2 shows the computed values of skin friction coefficient f''(0), couple stress -h'(0) and local Nusselt number  $-\theta'(0)$  for Cu- $Al_2O_3$ /water with different values of suction/injection parameter S and velocity slip parameter  $\alpha$ . It is found that f''(0), -h'(0) and  $-\theta'(0)$  are decreases with increasing S. The effect of Biot number,  $\gamma$  for different values of suction/injection parameter is given in Table 3. It shows that the heat transfer rate is a increasing function of  $\gamma$ . Numerical values of skin friction coefficient f''(0), and  $-\theta'(0)$  with different values of micropolar parameter K is given in table 4 and it is clear from this table, the values of f''(0), -h'(0), and  $-\theta'(0)$  increases with increasing K.

Table 1			
Thermophysical properties of base fluid and nanoparticles			
Physical properties	Fluid phase (water)	Cu	$Al_2O_3$
$C_{ ho}(J/kgK)$	4179	385	765
$\rho(kg/m^3)$	997.1	8933	3970
<i>κ</i> (W/mK)	0.613	400	40
$\sigma$ (S/m)	5.5 x 10 <sup>-6</sup>	5.9 x 10 <sup>6</sup>	35 x 10 <sup>6</sup>

Figures 2 to 4 show the variations of skin friction coefficient, couple stress coefficient and local Nusselt number with *S* for various values of velocity slip,  $\alpha$ . All variations show that it is decreases with *S* and increases with  $\alpha$ . The variations of local Nusselt number with *S* for various values of Biot number,  $\gamma$  is depicted in Figure 5. The variation is found increase with an increase in *S*. Temperature distribution which results in enhancement of local Nusselt number is due to physically increment both values of slip parameter,  $\alpha$  and Biot number,  $\gamma$ . Profiles for velocity, angular velocity and temperature distribution are plotted in Figures 6 to 8 for different values of velocity profile and temperature profile shown in Figures 7 and 8 increases with increase  $\alpha$ . This is due to increasing the velocity slip parameter and the portrayed that the boundary layer thickness also increase. Figure 9 is the plot of temperature profile for various values of Biot number,  $\gamma$  and it is shows that the temperature profile increases with increase  $\gamma$ . The increment shows in the temperature profile in Figures 8 and 9 proved that the increases in both parameters.



#### Table 2

Numerical values of of f''(0), -h'(0), and  $-\theta'(0)$  for Cu-A $l_2O_3$ /water with different values of  $\alpha$  when  $A^* = B^* = \theta_w = 1.2$ , n = 0.1,  $\gamma = 0.1$ ,  $\phi = 0.01$ ,  $M = 2.0 \times 10^{-12}$  Nr = 0.5 K = 0.1 Pr = 6.8

$0.01, M = 2.0 \times 10^{-12}, Nr = 0.5, K = 0.1, Pr = 6.8$					
S	α	f''(0)	-h'(0)	- heta'(0)	
-0.1651	0.4	0.4260	0.1661	0.1020	
-0.0651		0.2857	0.1481	0.1022	
0.0349		0.1575	0.1340	0.1020	
0.1349		0.0400	0.1234	0.1018	
0.2349		-0.0678	0.1159	0.1015	
-0.1939	0.5	0.4822	0.1766	0.1019	
-0.0939		0.3302	0.1558	0.1022	
0.0061		0.1944	0.1394	0.1021	
0.1061		0.0725	0.1267	0.1019	
0.2061		-0.0372	0.1175	0.1017	
-0.1861	0.6	0.4830	0.1797	0.1021	
-0.0861		0.3242	0.1573	0.1023	
0.0139		0.1854	0.1398	0.1021	
0.1139		0.0632	0.1264	0.1019	
0.2139		-0.0447	0.1165	0.1016	

#### Table 3

Numerical values of  $-\theta'(0)$  with different values  $\gamma$  when  $A^* = B^* = \theta_w = 1.2$ , n = 0.1,  $\phi = 0.01$ , M = 2.0 x  $10^{-12}$ , Nr = 0.5,  $\alpha = 0.1$ , K = 0.1, Pr = 6.8

S	γ	- heta'(0)
-0.1059	0.4	0.3682
-0.0059		0.3830
0.0941		0.3902
0.1941		0.3939
0.2941		0.3959
-0.1059	0.5	0.4457
-0.0059		0.4691
0.0941		0.4809
0.1941		0.4872
0.2941		0.4909
-0.1059	0.6	0.5185
-0.0059		0.5517
0.0941		0.5691
0.1941		0.5787
0.2941		0.5843

The variations of skin friction coefficient, couple stress coefficient and local Nusselt number with S for various values of material/micropolar parameter, K are plotted in Figures 10 to 12. All variations show that it is increases with S and increases with K. Physically, this is due to the material parameter supports the particles of the skew-symmetric of the fluid and therefore the couple stress coefficient increases. Figures 13 to 15 show the velocity profile, angular velocity profile and temperature profile for various values of material/micropolar parameter, K, respectively. As observed in Figures 13 and 14, the velocity profile and the angular velocity profile increases with increasing of micropolar parameter, K. Physically, this enhancement in the thickness of the momentum boundary layer indicates that the micropolar parameter produces less drag force in the boundary layer separation. From Figure 15, it is observed that the temperature profile decreases with increase material/micropolar parameter, K.



#### Table 4

Numerical values of $-\theta'(0)$	with different values $\gamma$	when $A^* = B^*$	$= \theta_w = 1.2$ ,
	10-12		

$n$ = 0.1, $\phi$ = 0.01, M = 2.0 x 10 <sup>-12</sup> , Nr = 0.5, $\alpha$ = 0.1, K = 0.1, Pr = 6.8					
S	Κ	f''(0)	-h'(0)	- heta'(0)	
0.2500	0.2	-0.0144	0.3166	0.1017	
0.3500		-0.1235	0.2958	0.1015	
0.4500		-0.2289	0.2795	0.1013	
0.5500		-0.3307	0.2675	0.1011	
0.6500		-0.4288	0.2598	0.1010	
0.1959	0.3	0.2265	0.7691	0.1020	
0.2959		0.1027	0.7028	0.1018	
0.3959		-0.0173	0.6418	0.1016	
0.4959		-0.1333	0.5865	0.1014	
0.5959		-0.2451	0.5376	0.1012	
0.0279	0.4	0.7827	1.7433	0.1026	
0.1279		0.6360	1.6140	0.1024	
0.2279		0.4929	1.4900	0.1022	
0.3279		0.3534	1.3714	0.1019	
0.4279		0.2174	1.2582	0.1017	



**Fig. 2.** Skin friction coefficient with *S* for different values of  $\alpha$ 



**Fig. 4.** Local Nusselt number with *S* for different values of  $\alpha$ 



Fig. 3. Couple stress with S for different values of  $\alpha$ 



**Fig. 5.** Local Nusselt number with *S* for different values of  $\gamma$ 





**Fig. 6.** Velocity profile for several values of  $\alpha$ 



Fig. 8. Temperature profile for several values of  $\alpha$ 



**Fig. 10.** Skin friction coefficient with *S* for different values of *K* 



**Fig. 7.** Angular velocity profile for several values of  $\alpha$ 



Fig. 9. Temperature profile for several values of  $\gamma$ 



**Fig. 11.** Couple stress with *S* for different values of *K* 





**Fig. 12.** Local Nusselt number with *S* for different values of *K* 



**Fig. 14.** Angular velocity profile for several values of *K* 



**Fig. 13.** Velocity profile for several values of *K* 



**Fig. 15.** Temperature profile for several values of *K* 

# 4. Conclusions

In the present work, the effects of velocity slip on MHD micropolar hybrid nanofluid over a stretching sheet with thermal radiation and non-uniform heat source/sink have been studied. The resulting partial differential equations, which describe the problem, are transformed into ordinary differential equations by using similarity transformations and the shooting method was implemented to get the solutions. Profiles for velocity, angular velocity and temperature are obtained for different values of the governing parameter involved. The results show that the skin friction and couple stress increase as the slip parameter,  $\alpha$  increase. This bring to the increases the thickness of boundary layer as increasing slip parameter,  $\alpha$ . Furthermore, the heat transfer coefficient increases as the velocity slip parameter,  $\alpha$  and biot number,  $\gamma$  increases. The angular velocity of the micropolar fluid increased for the higher values of material and slip parameters.



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